

Complex quantum states

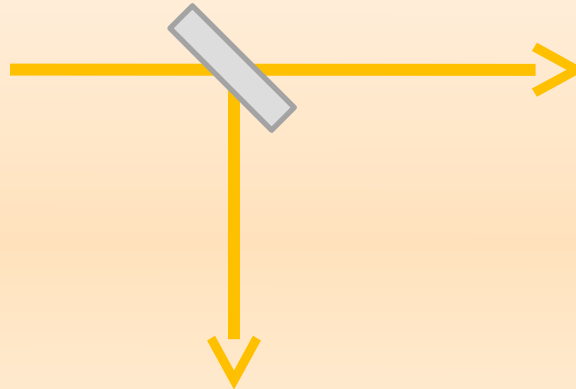
**Cai Yu, Wu Xingyao, Rafael Rabelo,
Le Huy Nguyen, Valerio Scarani**

motivation

Schrödinger cat?

“Superposition of two macroscopically distinguishable states”

Example:
A single photon on a beam-splitter



- “To be transmitted” and “to be reflected” are classical alternatives
- There are detectors that can distinguish those alternatives
- By arranging an interferometer, we know that the beam-splitter does create superposition

But this is not exactly what you had in mind, right?

Bigger kittens

Quantum optics: $|+\alpha\rangle + |-\alpha\rangle$ with α “large”

Monomode; not really achieved yet

Heavy weights: $\psi_1 + \psi_2$ for massive objects

Big molecules, nanomechanical oscillators,
current in SQUIDS...[*]

Also single degree of freedom

Qubit lover $|0\rangle^{\otimes N} + |1\rangle^{\otimes N}$

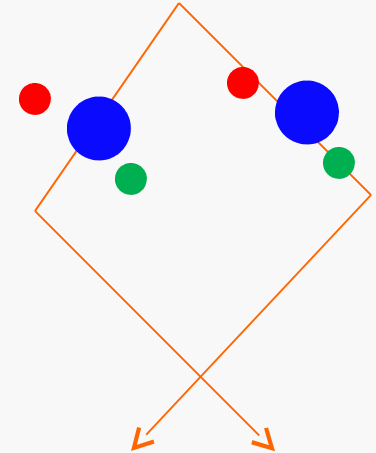
Many degrees of freedom, lots of
entanglement... but very easy to describe

[*] “we model the cat as a homogeneous sphere of water with a mass of 4kg”
Nimmrichter and Hornberger, PRL 2013, last sentence of Supplementary Information

A frequent confusion

“Interference of large molecules = GHZ state”, because

$$\begin{aligned}\psi_R + \psi_L = & |\vec{x}_R\rangle |\vec{x}_R + \vec{r}_1\rangle |\vec{x}_R + \vec{r}_2\rangle \dots \\ & + |\vec{x}_L\rangle |\vec{x}_L + \vec{r}_1\rangle |\vec{x}_L + \vec{r}_2\rangle \dots\end{aligned}$$



Disproof:

1) Write in relative coordinates (valid degrees of freedom)

$$\psi_R + \psi_L = (|\vec{x}_R\rangle + |\vec{x}_L\rangle) |\vec{r}_1\rangle |\vec{r}_2\rangle \dots$$

2) Notice that full purity is not needed for CM interference:

$$\psi_R + \psi_L = (|\vec{x}_R\rangle + |\vec{x}_L\rangle) \otimes \rho$$

(Sorry for the sloppy notation)

Complexity 101



- States useful for quantum computing must be “complex”
 - Simple states are easy to simulate
 - Critics of QC base their skepticism on the very possibility of creating such states
- Biological systems are also complex (OK, and in a hot environment too)



The Schrödinger cat is probably complex



- “Complexity is not computable”
 - Kolmogorov complexity is indeed not; other measures are: see next
- “Complex states are not feasible with current technologies”
 - True. But aren’t you tired of stuff that is “feasible with current technologies”? Shouldn’t theorists look a bit further ahead?

first steps in the forest

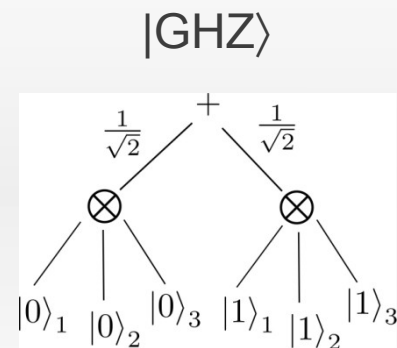
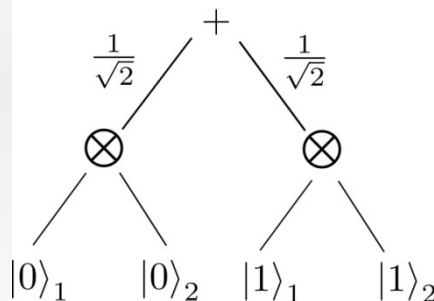
The “tree size” of a quantum state

Tree size of a quantum state

[Proposed by Aaronson STOC'04]

Any multiqubit quantum state can be described by a rooted tree of $\otimes$ and $+$ gates

Each leaf is labeled with a $\frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$ $|0\rangle + \beta|1\rangle$



- Size of a tree = **number of leaves.**
- **Tree size of a state (TS) = size of the *minimal* tree = most compact way of writing the state**

$$\underbrace{|0\rangle|0\rangle + |0\rangle|1\rangle + |1\rangle|0\rangle + |1\rangle|1\rangle}_{8 \text{ leaves}} = |+\rangle|+\rangle \quad \underbrace{|+\rangle|+\rangle}_{2 \text{ leaves}}$$

Getting compact (example: 3 qubits)

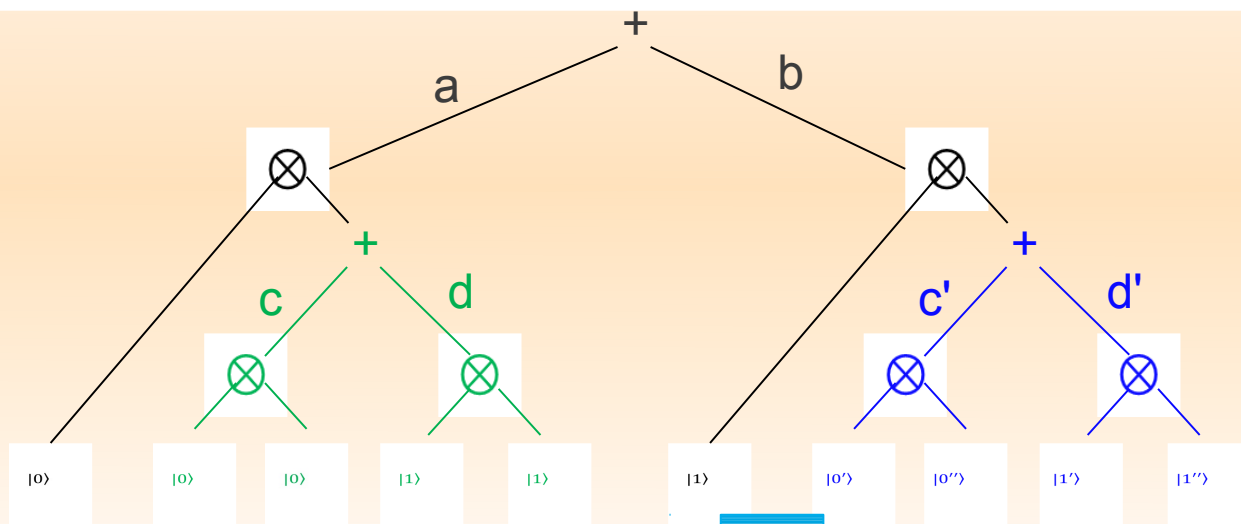
Naïve computational basis

$$\sum_{x_i=0,1} c_{x_1 x_2 x_3} |x_1\rangle |x_2\rangle |x_3\rangle$$

24 leaves

Sequential Schmidt:

$$a|0\rangle(c|0\rangle|0\rangle + d|1\rangle|1\rangle) + b|0\rangle(c'|0'\rangle|0''\rangle + d'|1'\rangle|1''\rangle)$$



10 leaves

Using Acin et al. 2001:

$$a|0\rangle|0\rangle|0\rangle + b|1\rangle(c|0'\rangle|0''\rangle + d|1'\rangle|1''\rangle)$$

8 leaves

Examples and bounds

$$TS(|0\rangle^{\otimes n}) = n$$

$$TS(GHZ) = 2n \text{ (I told you it's simple)}$$

$$TS(W) = O(n^2)$$

$$TS(1D \text{ cluster}) = O(n^4)$$

$O(n^2)$ this Monday's
result 😊

$$TS(2D \text{ cluster}) = 2^{\Omega(n)} \text{ conjectured}$$

$$TS(Shor) = n^{\Omega(\log n)} \text{ proved under one conjecture}$$

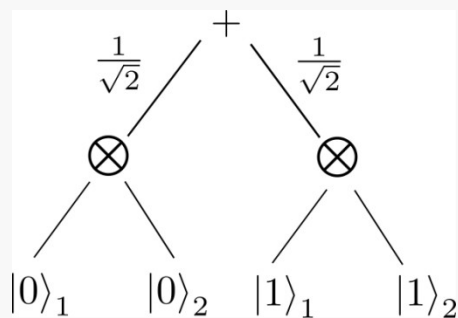
Anything provable
somewhere here?

- Upper bound on TS: nested Schmidt decomposition (or slightly more clever one) $\Rightarrow$ easy to prove that some states are NOT complex
- Conjectures: basically, if it allows universal QC, it cannot be too simple, otherwise we could simulate it.
- **Lower** bound on TS: size of a multilinear formula: see next

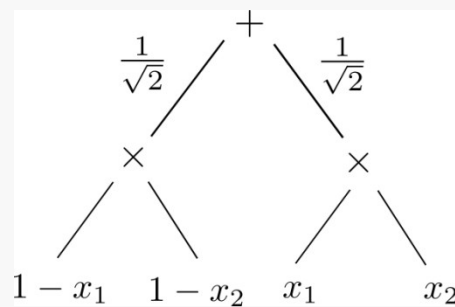
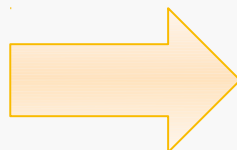
Multilinear formula

$$\Psi = \sum_{x_j=0,1} f(x_1, \dots, x_n) |x_1, \dots, x_n\rangle$$

We want a multilinear formula to compute the coefficients



$$\frac{1}{\sqrt{2}} (|00\rangle + |11\rangle)$$



$$\frac{1}{\sqrt{2}} [(1-x_1)(1-x_2) + x_1x_2]$$

So, what have we gained?

We can stand on the shoulder of mathematicians!

Raz, STOC'04: any multilinear formula that computes the **determinant** or **permanent** of a matrix is **super-polynomial**. So...

States with super-polynomial TS

arXiv:1303.4843

Take $n=m^2$ qubits, and arrange the coefficients as a matrix

$$\{x\} = \begin{pmatrix} x_{11} & x_{12} & \cdots & x_{1m} \\ x_{21} & x_{22} & \cdots & x_{2m} \\ \vdots & \vdots & \ddots & \vdots \\ x_{m1} & x_{m2} & \cdots & x_{mm} \end{pmatrix}.$$

Then these states have super-polynomial TS:

$$|\det_m\rangle = \sum_{x=0}^{2^n-1} \det(\{x\}) |x\rangle,$$

$$|\text{per}_m\rangle = \sum_{x=0}^{2^n-1} \text{perm}(\{x\}) |x\rangle$$

Raz

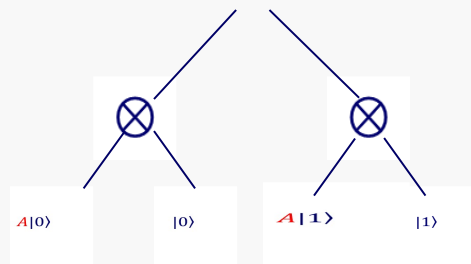
Our best
decomposition

$$n^{\Omega(\log n)/4} \leq TS \leq O((\sqrt{n})!)$$

Similar to the bound
“proved” for Shor

Work in progress (1): tight bounds

Tools: one can use the SLOCC classification, since the tree size does not change for a reversible LOCC:



3 qubits

- Biseparable: $TS = 5$
- GHZ class: $TS = 6$
- W class: $TS = 8$

Most complex, but of zero measure:

$$|W\rangle \rightarrow |W\rangle + \epsilon|111\rangle \in GHZ$$

4 qubits

- Several classes
- $TS \leq 14$ because any state can be written as $|0\rangle|GHZ\rangle + |1\rangle|GHZ'\rangle$ up to SLOCC
- There are states with $TS=14$ and their set is not of zero measure. E.g. $|0011\rangle + |0101\rangle + |1001\rangle + |0110\rangle + |1010\rangle + |1100\rangle$

Work in progress (2): TS & MPS

Assume for simplicity $n=2k$ qubits:

$$|MPS_d^n\rangle = \underbrace{(A_0^1|0\rangle + A_1^1|1\rangle)}_{1 \times d} \prod_{j=2}^{n-1} \underbrace{(A_0^j|0\rangle + A_1^j|1\rangle)}_{d \times d} \underbrace{(A_0^n|0\rangle + A_1^n|1\rangle)}_{d \times 1}$$

$$= \sum_{i=1}^D |MPS_{d,i}^{n/2}\rangle |MPS_{d,i}^{n/2}\rangle$$

Proof: insert $I_{d \times d} = e_1 e_1^T + \dots + e_d e_d^T$

$$\text{So } TS_k \leq 2d \times TS_{k-1} \leq \dots \leq (2d)^k$$

$$\text{that is } TS_{n,d} \leq (2d)^{\log_2 n} = n^{\log_2 2d}$$

- If d constant, TS polynomial in n $\square$
- 1D cluster state: $d=2$, so $TS=O(n^2)$.

Missing links

Wanted: operational interpretation(s)

- Links with measures of “difficulty”
 - Some polynomial circuit can generate superpolynomial TS states with finite probability
 - Rk: Shor’s algorithm also has polynomial circuit size...
 - No link known with Hamiltonian families (beyond indirect ones through MPS, see above)
- Link with universal quantum computing?
 - Clearly small TS $\Rightarrow$ easy to describe $\Rightarrow$ not useful
- Does high TS mean “harder to produce in the lab” even for few qubits?
- “Natural” situations in which such states appear
 - Biological systems??

Trying to enthuse Jose-Ignacio



Cai Yu (蔡宇)
PhD 3rd year

Rafael Rabelo
PhD 4th year

Lê Huy Nguyễn
Postdoc

Wu Xingyao
(邬兴尧)
PhD 1st year