

Bell inequalities for maximally entangled states

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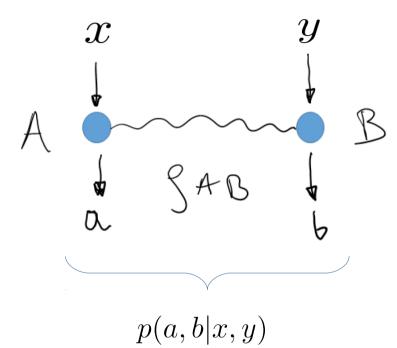


Foundation for Polish Science European Union European Regional Development Fund



Preliminaries

Bell scenario: consider two parties performing measurements on their local systems



Correlations are described by a collection of prob. distributions

 $\{p(a, b | x, y)\}_{a, b; x, y}$ $p(a, b | x, y) = \operatorname{Tr}[\rho_{AB}(M_x^a \otimes N_y^b)]$

measurement choices $x, y = 1, \dots, m$

(2,m,d) scenario

outcomes $a,b=0,\ldots,d-1$

measurement $M=\{M_a\}_a$ $M^a\geq 0, \quad \sum_a M^a=\mathbb{1}$ (POVM) $M^aM^{a'}=\delta_{a,a'}M^a$ (PM)

 $\{p(a,b|x,y)\}$ correlations/behavior

Preliminaries

Nonlocality and Bell inequalities

Local/classical correlations

$$p(a, b|x, y) = \sum_{\lambda} p(\lambda) p_{det}(a|x, \lambda) p_{det}(b|y, \lambda)$$

$$\forall_{a,b,x,y,\lambda} \quad p_{\det}(a|x,\lambda), p_{\det}(b|y,\lambda) \in \{0,1\}$$

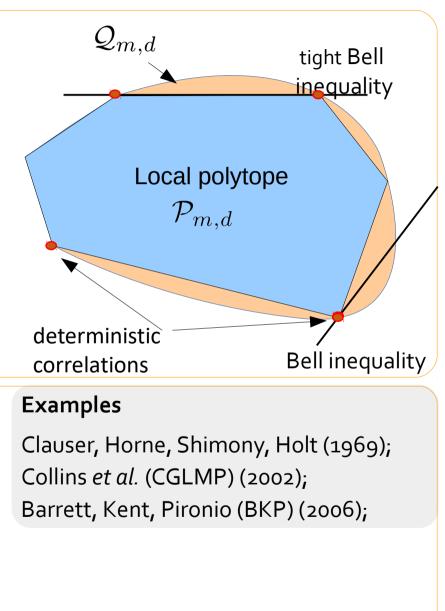
Otherwise they are called nonlocal (nonlocality)

 $\mathcal{P}_{m,d} \subsetneq \mathcal{Q}_{m,d}$ [J. S. Bell, Physics **1**, 195 (1964)]

 Bell inequalities: Hyperplanes constraining the local set

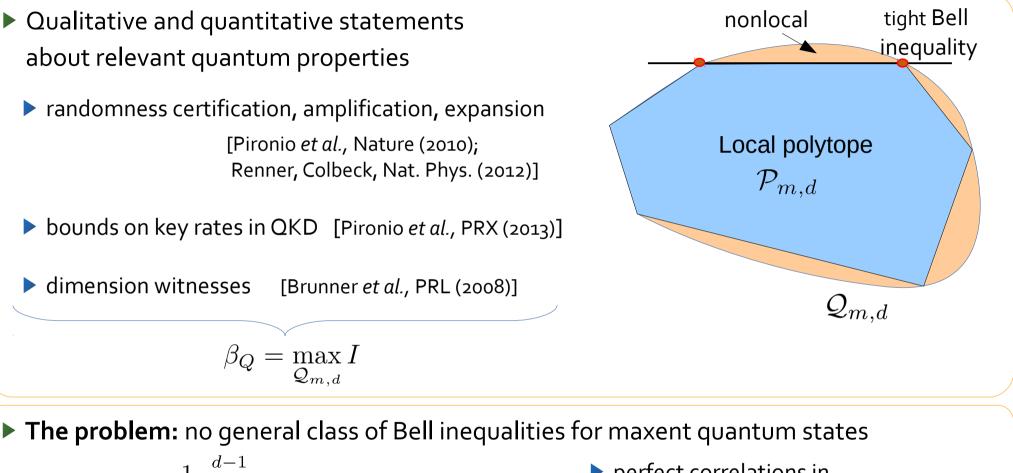
$$I := \sum_{a,b,x,y} T_{x,y}^{a,b} p(a,b|x,y) \le \beta_C$$
$$(\beta_C = \max_{\mathcal{P}_{m,d}} I$$

Finite number of BI's (facets) enough to fully characterize the local polytope



tight Bell inequalities (convex hull problem)

The problem



$$|\psi_d^+\rangle = \frac{1}{\sqrt{d}} \sum_{i=0}^{d-1} |i\rangle_A \otimes |i\rangle_B \in \mathbb{C}^d \otimes \mathbb{C}^d$$

- CHSH (1971) (2,2,2) scenario
- Son *et al.* (2006) (2,2,d) scenario
- Ji et al. (2008), Liang et al. (2009) (2,2,3) scenario
 MUB observables

- perfect correlations in any local basis (QKD)
- $\rho_A = \rho_B = \mathbb{1}/d$ (randomness certification)
- maximizer of entanglement measures

Constructing Bell inequalities

Phys. Rev. Lett. 119, 040402 (2017)

Consider the Barrett-Kent-Pironio (BKP) Bell expression

p(a, b|x, y)

[Masanes, QIC (2002)]

- facet Bell inequalities in (2,2,d) scenario
- \blacktriangleright not maximally violated by $|\psi_d^+
 angle$

 $\lfloor d/2 \rfloor {-}1$

 $I_{d,m} := \sum (\mathbb{P}_k - \mathbb{Q}_k)$

▶ e.g., for *d*=3

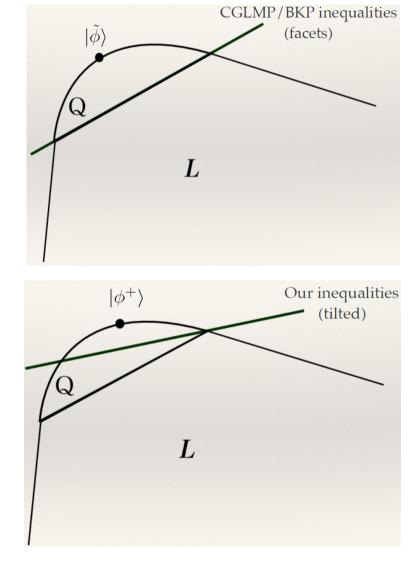
 $|\widetilde{\phi}\rangle \sim |00\rangle + \gamma|11\rangle + |22\rangle \qquad \gamma = \frac{1}{2}(\sqrt{11} - \sqrt{3})$

[Acin, Durt, Gisin, QIC (2002); Yang et al. (2014)]

Modify by adding parameters (tilting the inequality)

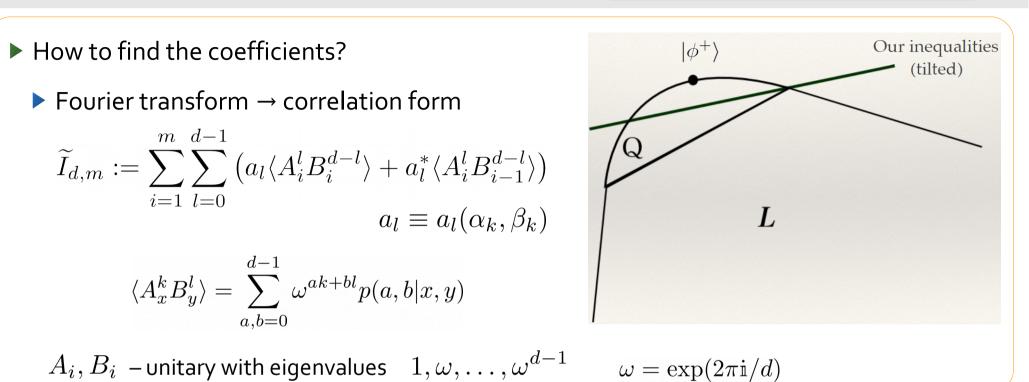
$$I_{d,m} := \sum_{k=0}^{\lfloor d/2 \rfloor - 1} \left(\alpha_k \mathbb{P}_k - \beta_k \mathbb{Q}_k \right)$$
$$\alpha_k, \beta_k \in \mathbb{R}$$

[Collins *et al.*, PRL (2002); Barrett *et al.*, PRL (2006)]



Constructing Bell inequalities

Phys. Rev. Lett. 119, 040402 (2017)



''Quantum approach'' (CHSH example)

$$I_{\text{CHSH}} = \langle A_0 \otimes (B_0 + B_1) \rangle + \langle A_1 \otimes (B_0 - B_1) \rangle \le 2$$

optimal CHSH measurements

$$A_0 \otimes \frac{B_0 + B_1}{\sqrt{2}} |\psi_2^+\rangle = \sigma_x \otimes \sigma_x |\psi_2^+\rangle = |\psi_2^+\rangle \qquad A_{0/1} = \sigma_{x/z}, \quad B_{0/1} = \frac{\sigma_x \pm \sigma_z}{\sqrt{2}}$$
$$A_1 \otimes \frac{B_0 - B_1}{\sqrt{2}} |\psi_2^+\rangle = \sigma_z \otimes \sigma_z |\psi_2^+\rangle = |\psi_2^+\rangle \qquad |\psi_2^+\rangle = \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle)$$

Constructing Bell inequalities

"Quantum approach"

$$\widetilde{I}_{d,m} := \sum_{i=1}^{m} \sum_{l=1}^{d-1} \left(a_l \langle A_i^l B_i^{d-l} \rangle + a_l^* \langle A_i^l B_{i-1}^{d-l} \rangle \right) = \sum_{i=1}^{m} \sum_{l=1}^{d-1} \langle A_i^l \overline{B}_i^{(l)} \rangle$$
$$a_l \equiv a_l(\alpha_k, \beta_k) \qquad \overline{B}_i^{(l)} = a_l B_i^{d-l} + a_l^* B_{i-1}^{d-l}$$

• we want α_k, β_k so that

Full characterization of our Bell inequalities Phys. Rev. Lett. 119, 040402 (2017) Analytical proof of the maximal quantum value $\beta_Q = \max_{\mathcal{Q}_{m,d}} \tilde{I}_{m,d} = m(d-1)$ Local polytope sum of squares decomposition $\mathcal{P}_{m,d}$ $\beta_Q \mathbb{1} - \mathcal{B} = \frac{1}{2} \sum_{k=1}^{m} \sum_{k=1}^{d-1} P_{ik}^{\dagger} P_{ik} + \frac{1}{2} \sum_{k=1}^{m-2} \sum_{k=1}^{d-1} T_{ik}^{\dagger} T_{ik}$ $\mathbb{Q}_{m,d}$ $P_{ik} = \mathbb{1} - A_i^k \otimes \bar{B}_i^k \qquad T_{ik} = \mu_{i,k} B_2^{d-k} + \nu_{i,k} B_{i+2}^{d-k} + \tau_{ik} B_{i+3}^{d-k}$ ▼ ↑ ▼ some coefficients

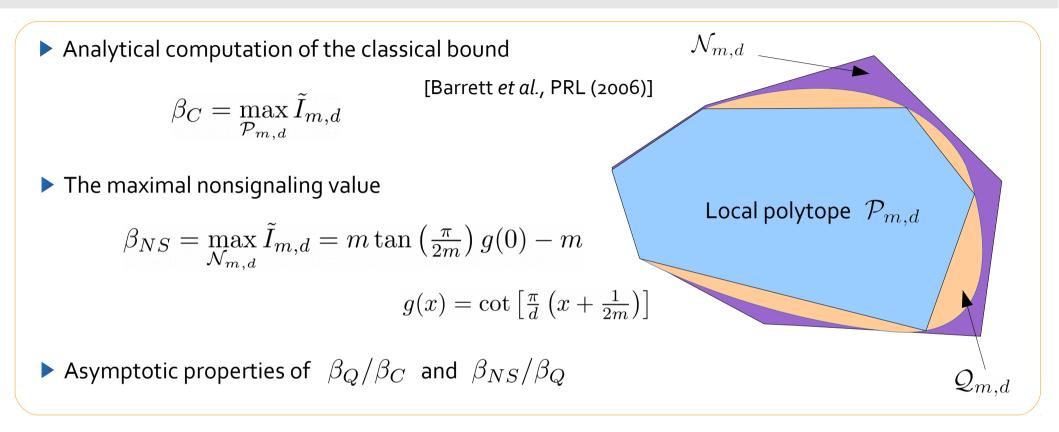
• quantum realization $|\psi_d^+
angle$ and optimal CGLMP/BKP measurements

Example: the CHSH inequality

$$2\sqrt{2}\,\mathbb{1} - \mathcal{B}_{\text{CHSH}} = \frac{1}{\sqrt{2}} \left[\left(\mathbb{1} - A_0 \otimes \frac{B_0 + B_1}{\sqrt{2}}\right)^2 + \left(\mathbb{1} - A_1 \otimes \frac{B_0 - B_1}{\sqrt{2}}\right)^2 \right]$$

for all dichotomic A_i, B_i

Full characterization of our Bell inequalities



Special cases

m=2 – Bell inequalities considered by Son, Lee and Kim

[Phys. Rev. Lett. (2006); J. de Vicente, Phys. Rev. A (2015)]

d=2 – the chained Bell inequalities

[Pearle, 1970; Braunstein, C. Caves, 1990; Wehner (2006)]

$$\langle A_0 B_0 \rangle + \langle A_0 B_m \rangle + \sum_{i=1}^m \left(\langle A_i B_i \rangle + \langle A_i B_{i-1} \rangle \right) \le 2(m-1)$$

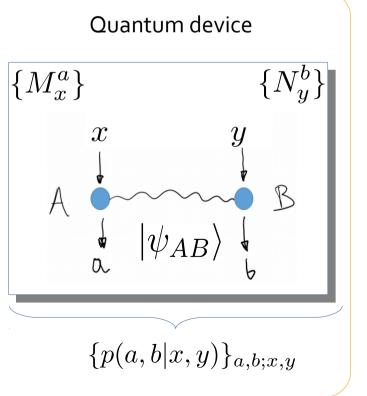
Applications

- Device-independent applications
 - Quantum key distribution better lower bounds on key rates
 - Self-testing of maximally entangled states for d>2

- The idea of self-testing (DI certification)
 - $\{p(a,b|x,y)\}_{a,b;x,y}$
 - or violation of some Bell inequality

$$\sum_{a,b,x,y} T^{a,b}_{x,y} \ p(a,b|x,y) = \beta$$

• deduce properties of $|\psi_{AB}\rangle$ and $\{M_x^a\}, \{N_y^b\}$

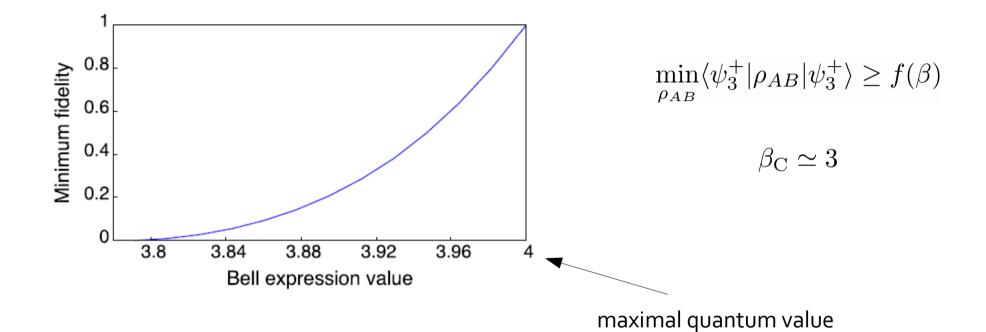


Self-testing maxent state

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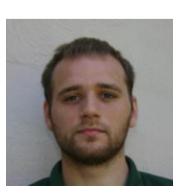
Numerical for d=3 (m=2)

[Yang et al., 2014; Wu et al., 2014]



Rigorous analytical proof for d=3,4,5

see Jędrek's talk!



An approach based on the CHSH Bell inequality [Yang, Navascués, 2012]

Generalization to N parties

Bell inequalities maximally violated by N qudit GHZ states

$$|\mathrm{GHZ}_{N,d}\rangle = \frac{1}{\sqrt{d}} \sum_{i=0}^{d-1} |i\rangle^{\otimes N}$$

Consider the Bell expression (N=3) [Aolita et al., PRL (2012); Bancal et al. JPA (2012)]

$$I_{N,d,m} := \sum_{k=0}^{\lfloor d/2 \rfloor - 1} (\alpha_k \mathbb{P}_k - \beta_k \mathbb{Q}_k) \qquad \qquad \alpha_k, \beta_k \in \mathbb{R} \quad \begin{array}{c} \text{free} \\ \text{parameters} \end{array}$$

$$p(a, b, c | x, y, z) \qquad \qquad \alpha_k = \beta_k = 1 - \frac{2k}{d-1}$$

Go to the correlation picture (the Fourier transform)

$$\tilde{I}_{3,M,d} = \frac{1}{d} \sum_{x,y=1}^{M} \sum_{k=0}^{d-1} \langle \bar{A}_x^k B_{x+y-1}^{-k} C_y^k \rangle,$$

 A_x, B_y, C_z – unitary observables with eigenvalues $1, \omega, \dots, \omega^{d-1}$

$$\bar{A}_x^k = a_k(\vec{\alpha}, \vec{\beta})A_x^k + a_k^*(\vec{\alpha}, \vec{\beta})A_{x+1}^k$$

Generalization to N parties

Solve a system of equations

$$\bar{A}_x^k \otimes B_{x+y-1}^{-k} \otimes C_y^k | \text{GHZ}_{3,d} \rangle = | \text{GHZ}_{3,d} \rangle \qquad \text{for all } x, y, k$$

for fixed A_x, B_y, C_z α_k, β_k almost uniquely

Characterization

Maximal quantum violation (SOS similar to the bipartite one)

 $\beta_{N,m,d}^Q = m^{N-1}(d-1) = m^{N-2}\beta_{2,m,d}^Q$

The maximal nonsignaling value

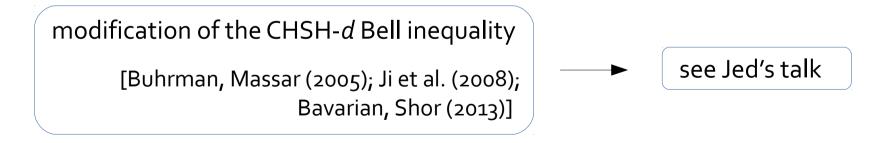
 $\beta_{N,m,d}^{\rm NS} = m^{N-2} \beta_{2,m,d}^{\rm NS}$

Maximal classical value

 $\beta_{N,m,d}^{\mathcal{C}} = ? \qquad \beta_{N,m,d}^{\mathcal{C}} \neq m^{N-2} \beta_{2,m,d}^{\mathcal{C}}$

Conclusion/Outlook

- A class of Bell inequalities for maxent states
- Analytical computation of all relevant quantities $\beta_C, \beta_Q, \beta_{NS}$
- ► Numerical self-testing for *d*=3 —
- Generalization to the GHZ states
- Self-testing of maxent states for any d?
- Other Bell inequalities max. violated by the max. entangled states?



Bell inequalities for other entangled quantum states of local dimension >2

partially entangled bipartite states

multipartite states such as AME states or graph states

for the rigorous analytical

proof see Jed's talk